# Computational Efficient Parameter and Performance Prediction of Wound-Rotor Induction Motor

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A basic, fast, and accurate steady-state finite-element analysis (FEA) based parameter and performance calculation method for the wound-rotor induction motor is presented. The method couples the circuit equations with field equations using the principle of frozen permeabilities to present the voltage-fed induction motor. The effects of saturation saliency and cross-magnetization are accurately taken into account. The method is verified by direct comparison with the current- and torque-slip frequency curves obtained from a commercial FEA package.

Index Terms—Cross-coupling, finite-element analysis (FEA), flux linkages, two-axis modeling, wound-rotor induction motor (WR-IM).

## I. INTRODUCTION

ARGE power wound-rotor induction motors (WR-IMs) are still one of the common motor loads on the grid, widely used for high locked rotor torque and low startup currents [1]–[4]. Hence, fast and accurate parameter determination of the WR-IM is of topical interest; with these parameters, the performance characteristics can be determined for the optimum design of the motor [5]–[10].

The main drawback is the complex analysis of the WR-IM because of the existence of alternating currents in both stator and rotor of the motor. The finite-element (FE) analysis (FEA) method is a broadly used, very accurate numerical method for motor performance determination [11]–[16]. The time-stepping FEA method has been used for the accurate determination of the transient performance of the IM [3], [5], [17]–[19]. However, the method is not computationally time efficient for steady-state analysis and machine design optimization.

Lately, the time-harmonic (TH) FEA method, mainly applied to squirrel-cage induction motors, has been widely used for strict computation of the motor performance [18], [20]-[23]. The major drawback of the stated TH method is that only the linear region of the motor is considered and that is also time consuming especially when only steady-state performance is needed. Mezani et al. [16] successfully use the truncated version of the TH method [23], [24] for the first time on the computation of the WR-IM. In [16], the double air-gap method is used, only considering the principal airgap space-harmonic of the magnetic flux in which only four FE magnetostatic computations are required to determine the WR-IM performance for any slip value. The disadvantage of this method is that only the linear case or in average non-linear case is considered [23]. The accuracy of the proposed method of Mezani et al. [16] has proved not always to be, essentially on the torque prediction.

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Instead of considering the first space-harmonic air-gap magnetic flux proposed in [16], the total overall flux linkage of the motor produced by the three phases is considered in this paper for parameter and performance calculation. A simple and basic FE parameter and performance prediction method recommended for the design optimization of WR-IM is newly introduced. With this method not only the linear case but also the saturation saliency and cross-coupling effects on the motor parameters incorporating all leakage flux linkages are considered. A proof is given on the extreme fastness and accuracy of the method in performance prediction.

#### II. FLUX LINKAGE COMPUTATION OF WR-IM

In the total flux linkage computation, the field is assumed not changing in the *z*-axis. Thus, the 2D FE solution magnetic vector potential has only the *z*-directed component, i.e.,

$$\mathbf{A} = A_z \mathbf{z} \tag{1}$$

where **z** is the *z*-directed unit vector. The net total flux linkage of a winding through a surface *S* is calculated using the relation  $\mathbf{B} = \nabla \times \mathbf{A}$  and Stoke's theorem as

$$\boldsymbol{\varphi} = \int_{S} \mathbf{B} \cdot d\mathbf{S} = \int_{S} (\boldsymbol{\nabla} \times \mathbf{A}) \cdot d\mathbf{S} = \oint_{C} \mathbf{A} \cdot d\mathbf{I}.$$
 (2)

Considering (2) for a coil area S divided into n first-order triangle mesh FEs, N turns and length l along the z-direction, the flux linkage of a coil is

$$\lambda = N\boldsymbol{\varphi} = N\sum_{j=1}^{n} \frac{\boldsymbol{\Delta}_{j}}{S} \left(\Im \frac{1}{3} \sum_{i=1}^{3} \mathbf{A}_{ij}\right) l.$$
(3)

In (3),  $\mathbf{A}_{ij}$  is the magnetic vector potential of nodal point i = 1, i = 2, or i = 3 of the *j*th triangular element of area  $\mathbf{\Delta}_j$ . The direction of the integral either into or out of a plane is given by  $\Im$  (i.e., +1 or -1). Considering (3) and one meshed pole, the net total flux linkage of a phase winding can be expressed as

$$\lambda_{\rm abc} = \frac{2pNl}{S} \sum_{j=1}^{u} \left( \frac{\mathbf{\Delta}_j \Im}{\Im} \sum_{i=1}^{3} \mathbf{A}_{ij} \right) \tag{4}$$

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where u is the total number of elements of the meshed coil areas of the phase in the pole region [25]. The total net flux linkage of (4) includes high-order harmonic and leakage fluxes produced by the winding, slotted air gap, and magnetic saturation. The total flux linkage is essential because with that a complete parameter and performance (e.g., power factor) prediction model of the motor can be obtained.

## III. ELECTROMAGNETIC MODEL OF WR-IM

The stator and rotor voltages are classically expressed as

$$\mathbf{v}_{\mathrm{abc}_{(s,r)}} = \mathbf{R}_{(s,r)}\mathbf{i}_{\mathrm{abc}_{(s,r)}} + \frac{d}{dt}\boldsymbol{\lambda}_{\mathrm{abc}_{(s,r)}}$$
(5)

where  $R_{(s,r)}$  given by

$$R_{(s,r)} = (\rho_{cu})_t \frac{l_{T_{(s,r)}} N_{(s,r)}^2 pq}{A_{cu_{(s,r)}}} \qquad l_{T_{(s,r)}} = 2(l + l_{e_{(s,r)}}) \tag{6}$$

and  $\lambda_{abc_{(s,r)}}$  of (4) are stator or rotor phase resistances and flux linkages. In (6),  $A_{cu_{(s,r)}}$  and  $(\rho_{cu})_T$  are the copper slot crosssectional area and resistivity at temperature *T*, respectively. The total conductor length  $l_T$  of (6) is the sum of stack length *l* and end-winding length  $l_e$  given by Boldea and Nasar [11].  $N_{(s,r)}$  and *q* are the number of turns per slot and slots per pole per phase, respectively. The balanced three voltage, current, and flux linkage phases of (5), representing a stationary time-dependent system, are transformed to a rotating time-dependent system (DQ-axis) and inverse transformed as

$$\mathbf{v}_{dq_{(s,r)}} = \mathbf{T}(\theta_{(s,r)})\mathbf{v}_{abc_{(s,r)}}; \quad \mathbf{v}_{abc_{(s,r)}} = \mathbf{T}^{-1}(\theta_{(s,r)})\mathbf{v}_{dq_{(s,r)}}$$
(7)

$$\mathbf{i}_{dq(s,r)} = \mathbf{T}(\theta_{(s,r)})\mathbf{i}_{abc_{(s,r)}}; \quad \mathbf{i}_{abc_{(s,r)}} = \mathbf{T}^{-1}(\theta_{(s,r)})\mathbf{i}_{dq(s,r)}$$
(8)

$$\boldsymbol{\lambda}_{dq_{(s,r)}} = \mathbf{T}(\boldsymbol{\theta}_{(s,r)})\boldsymbol{\lambda}_{\mathrm{abc}_{(s,r)}}; \quad \boldsymbol{\lambda}_{\mathrm{abc}_{(s,r)}} = \mathbf{T}^{-1}(\boldsymbol{\theta}_{(s,r)})\boldsymbol{\lambda}_{dq_{(s,r)}} \quad (9)$$

where **T** and  $\mathbf{T}^{-1}$  are transformation matrices expressed in terms of the electrical rotor position angle (transformation angle)  $\theta$  given in matrix form as

$$\mathbf{T} = \begin{bmatrix} \cos\theta & \cos(\theta - 120^{\circ}) & \cos(\theta + 120^{\circ}) \\ -\sin\theta & -\sin(\theta - 120^{\circ}) & -\sin(\theta + 120^{\circ}) \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}.$$
 (10)

The inverse transformation is expressed as

$$\mathbf{T}^{-1} = \begin{bmatrix} \cos\theta & -\sin\theta & 1\\ \cos(\theta - 120^\circ) & -\sin(\theta - 120^\circ) & 1\\ \cos(\theta + 120^\circ) & -\sin(\theta + 120^\circ) & 1 \end{bmatrix}.$$
 (11)

For the stator and rotor of the WR-IM,  $\theta$  of (10) and (11) is replaced by

$$\theta_s = \omega_s t + \theta_0 = 2\pi f_s t + \theta_0$$
  
$$\theta_r = (\omega_s - \omega_r)t + \theta_0 = \omega t + \theta_0 = 2\pi (f_s - f_r)t + \theta_0 \qquad (12)$$

respectively, where t and  $\theta_0$  are the rotor position time and the initial rotor position, respectively.  $\omega_s$  and  $\omega_r$  are the stator and rotor angular velocities, respectively, and  $f_s$  and  $f_r$  are the grid supply and rotor frequencies, respectively. The slip frequency f of (12) for a p-pole pair motor is calculated from the rotor speed  $n_r$  as

$$f = (f_s - f_r) = \left(f_s - \frac{n_r p}{60}\right).$$
 (13)



Fig. 1. (a) D- and (b) Q-axes equivalent circuit in the synchronously rotating reference frame in steady state of the WR-IM.



Fig. 2. DQ-axis phasor diagram of the WR-IM.

The transformed **T** equations of (5) in (7)–(9) result in

$$\mathbf{v}_{dqs} = \mathbf{R}_{s} \mathbf{i}_{dqs} + \frac{d}{dt} \boldsymbol{\lambda}_{dqs} \mp \boldsymbol{\omega}_{s} \boldsymbol{\lambda}_{qds}$$
$$\mathbf{v}_{dqr} = \mathbf{R}_{r} \mathbf{i}_{dqr} + \frac{d}{dt} \boldsymbol{\lambda}_{dqr} \mp (\boldsymbol{\omega}_{s} - \boldsymbol{\omega}_{r}) \boldsymbol{\lambda}_{qdr}.$$
(14)

In steady state, the transformed voltage equations of (14) become

$$\mathbf{V}_{dqs} = \mathbf{R}_{s} \mathbf{I}_{dqs} \mp \boldsymbol{\omega}_{s} \boldsymbol{\lambda}_{qds}$$
$$\mathbf{V}_{dqr} = \mathbf{R}_{r} \mathbf{I}_{dqr} \mp (\boldsymbol{\omega}_{s} - \boldsymbol{\omega}_{r}) \boldsymbol{\lambda}_{qdr}.$$
(15)

The voltage equations of (15), including the effect of the endwinding inductance  $L_e$ , which is given by

$$L_{e_{(s,r)}} = V_{(u)} m D_{(s,r)} (N_{(s,r)} q k_d k_{p(u)})^2 k_{e(p)} (10)^{-8}$$
(16)

where *m* is the number of phases and the shape factor  $V_{(u)}$ , distribution factor  $k_d$ , pitch factor  $k_p$ , and end-winding factor  $k_e$  explicitly given by Kamper [26], result in the DQ-axis equivalent circuits and phasor diagrams of Figs. 1 and 2, respectively. The core losses are not included in the equivalent circuit (by means of a core loss resistance) because these losses are lumped with mechanical losses as rotational losses.

The total DQ-axis flux linkages  $\lambda_{qd_{(s,r)}}$  of (15) can be segregated into components and represented in matrix form given by

$$\boldsymbol{\lambda}_{dq_{(s,r)}} = \mathbf{L}_{(s,r)} \mathbf{I}_{dq_{(s,r)}} \tag{17}$$

as

$$\begin{bmatrix} \lambda_{ds} \\ \lambda_{qs} \\ \lambda_{dr} \\ \lambda_{qr} \end{bmatrix} = \begin{bmatrix} L_{ds} & M_{dsqs} & M_{dsdr} & M_{dsqr} \\ M_{qsds} & L_{qs} & M_{qsdr} & M_{qsqr} \\ M_{drds} & M_{drqs} & L_{dr} & M_{drqr} \\ M_{qrds} & M_{qrqs} & M_{qrdr} & L_{qr} \end{bmatrix} \begin{bmatrix} I_{ds} \\ I_{qs} \\ I_{dr} \\ I_{qr} \end{bmatrix}$$
(18)

where L and M represent the self and mutual inductances of the motor windings, respectively. In (17),  $\mathbf{L}_{(s,r)}$  is the stator and rotor inductance matrix.

The performance equations of (15) can be modified by using (18) to become

$$\mathbf{V} = \mathbf{Z}\mathbf{I} \tag{19}$$

where V, I and Z are the voltage, current and impedance matrices given by

$$\mathbf{V} = \begin{bmatrix} I_{ds} & I_{qs} & I_{dr} & I_{qr} \end{bmatrix}$$
(20)

$$\mathbf{I} = \begin{bmatrix} I_{ds} & I_{qs} & I_{dr} & I_{qr} \end{bmatrix}$$
(21)

and (22), shown at the bottom of the page, respectively. It can be seen from the performance equation of (19) that knowing V and Z of (20) and (22), respectively, I of (21) can be calculated as

$$\mathbf{I} = \mathbf{Z}^{-1} \mathbf{V}.$$
 (23)

Also following the transformed flux linkages  $\lambda_{dq_{(s,r)}}$  of (9), the general average total torque of the WR-IM is given by

$$T_t = \frac{3}{2} p(\lambda_{ds} I_{qs} - \lambda_{qs} I_{ds}).$$
(24)

The total torque can further be segregated into components by using the flux linkage components of (18) in (17) into the total torque equation of (24) to give the total torque as

$$T_t = T_r + T_s + T_m \tag{25}$$

where

$$T_{r} = \frac{3}{2} p[(M_{dsqr}I_{qr} + M_{dsdr}I_{dr})I_{qs} - (M_{qsdr}I_{dr} + M_{qsqr}I_{qr})I_{ds}]$$

$$T_{s} = \frac{3}{2} p[(L_{ds} - L_{qs})I_{ds}I_{qs}]$$

$$T_{m} = \frac{3}{2} p[M_{dq}(I_{qs}^{2} - I_{ds}^{2})]\{M_{dsqs} = M_{qsds} = M_{dq}\}.$$
 (26)

In (25), the total torque  $T_t$  is expressed in terms of the rotor flux torque  $(T_r)$ , saliency torque  $(T_s)$ , and cross-magnetizing torque  $(T_m)$ .  $T_r$  is due to the rotor-induced flux seen by the stator,  $T_s$  is due to the difference between the two D- and Q-axes stator inductances engendering in reluctance or saliency torque and  $T_m$  is due to the cross-magnetization between the D- and Q-axes. With regard to calculation of the torque, it is of great importance to understand that the torque of (24) does not include higher order permeance harmonics torques, i.e., the flux derivative terms have been neglected in (24).

Equations (23) and (25) form the basis of parameter and performance prediction of the WR-IM of the method described in Sections IV and V.

## IV. FEA COMPUTATION OF WR-IM INDUCTANCES

To calculate the inductance matrix  $\mathbf{L}_{(s,r)}$  of (18) in (17), the total flux linkages  $\lambda_{dq_{(s,r)}}$  of (17) have to be segregated as in (18). To do so, the frozen permeability method [27]–[29], implemented in this paper for the WR-IM, is used to compute the flux linkage components.

In the method, a double excited non-linear FEA solution is used to determine the total transformed DQ-axis flux linkages  $\lambda_{dq(s,r)}$  of (18) in (17). Thus, an initial current  $\mathbf{I}_{dq(s,r)}$  of (8) is set by current angles  $\alpha_s$  and  $\alpha_r$  of Fig. 2. The  $\mathbf{I}_{dq(s,r)}$  is inverse transformed to obtain phase current  $\mathbf{i}_{abc(s,r)}$  as in (8), which is in turn used by the FEA package to calculate the phase flux linkage  $\lambda_{abc(s,r)}$ .  $\lambda_{abc(s,r)}$  is inversely transformed to obtain the total DQ-axis flux linkages  $\lambda_{dq(s,r)}$ . The permeabilities of the elements of the FEA mesh are then frozen to preserve all the information about saturation in the motor on the stated load condition. Thus, by freezing the permeabilities the nonlinear problem becomes a linear problem. In the instance of the WR-IM, four linear magnetostatic FEA solutions are necessary to calculate the DQ-axis flux linkage components.

The columns of the inductance matrix  $\mathbf{L}_{(s,r)}$  of (18) in (17) are defined by

$$\mathbf{L}_{(s,r)_{1}} = \begin{bmatrix} L_{ds} & M_{qsds} & M_{drds} & M_{qrds} \end{bmatrix}^{T}$$
$$\mathbf{L}_{(s,r)_{2}} = \begin{bmatrix} M_{dsqs} & L_{qs} & M_{drqs} & M_{qrqs} \end{bmatrix}^{T}$$
$$\mathbf{L}_{(s,r)_{3}} = \begin{bmatrix} M_{dsdr} & M_{qsdr} & L_{dr} & M_{qrdr} \end{bmatrix}^{T}$$
$$\mathbf{L}_{(s,r)_{4}} = \begin{bmatrix} M_{dsqr} & M_{qsqr} & M_{drqr} & L_{qr} \end{bmatrix}^{T}.$$
(27)

According to (18)

$$\lambda_{dq(s,r)_{1}} = \mathbf{L}_{(s,r)_{1}} I_{ds}$$

$$\lambda_{dq(s,r)_{2}} = \mathbf{L}_{(s,r)_{2}} I_{qs}$$

$$\lambda_{dq(s,r)_{3}} = \mathbf{L}_{(s,r)_{3}} I_{dr}$$

$$\lambda_{dq(s,r)_{4}} = \mathbf{L}_{(s,r)_{4}} I_{qr}.$$
(28)

The flux linkages of (28) are determined from the four linear magnetostatic FEA solutions with single excitation of the DQ-axis currents as follows:

$$\lambda_{dq_{(s,r)_{1}}} \qquad \mathbf{I} = \begin{bmatrix} I_{ds} & 0 & 0 & 0 \end{bmatrix}^{T}$$

$$\lambda_{dq_{(s,r)_{2}}} \qquad \mathbf{I} = \begin{bmatrix} 0 & I_{qs} & 0 & 0 \end{bmatrix}^{T}$$

$$\lambda_{dq_{(s,r)_{3}}} \qquad \mathbf{I} = \begin{bmatrix} 0 & 0 & I_{dr} & 0 \end{bmatrix}^{T}$$

$$\lambda_{dq_{(s,r)_{4}}} \qquad \mathbf{I} = \begin{bmatrix} 0 & 0 & 0 & I_{qr} \end{bmatrix}^{T}. \qquad (29)$$

$$\mathbf{Z} = \begin{bmatrix} (R_s - \omega_s M_{qsds}) & -\omega_s (L_{e_s} + L_{qs}) & -2\pi f_s M_{qsdr} & -2\pi f_s M_{qsqr} \\ \omega_s (L_{e_s} + L_{ds}) & (R_s + \omega_s M_{dsqs}) & 2\pi f_s M_{dsdr} & 2\pi f_s M_{dsqr} \\ -\omega M_{qrds} & -\omega M_{qrqs} & (R_r - \omega M_{qrdr}) & -\omega (L_{e_r} + L_{qr}) \\ \omega M_{drds} & \omega M_{drqs} & \omega (L_{e_r} + L_{dr}) & (R_r + \omega M_{drqr}) \end{bmatrix} \end{bmatrix}$$
(22)



Fig. 3. Parameter and performance prediction flow-diagram of the WR-IM.

With the flux linkages of (29) known, the inductances of (27) using (28) are given by

$$\mathbf{L}_{(s,r)_{1}} = \lambda_{dq_{(s,r)_{1}}} I_{ds}^{-1} 
\mathbf{L}_{(s,r)_{2}} = \lambda_{dq_{(s,r)_{2}}} I_{qs}^{-1} 
\mathbf{L}_{(s,r)_{3}} = \lambda_{dq_{(s,r)_{3}}} I_{dr}^{-1} 
\mathbf{L}_{(s,r)_{4}} = \lambda_{dq_{(s,r)_{4}}} I_{qr}^{-1}.$$
(30)

It is of great significance to take note that the total DQaxis non-linear FEA solution is equal to the sum of the four respective individual DQ-axis linear FEA solutions i.e.,

$$\lambda_{dq_{(s,r)}} = \sum_{i=1}^{4} \lambda_{dq_{(s,r)i}} = \sum_{i=1}^{4} \mathbf{T}(\theta_{(s,r)}) \lambda_{abc_{(s,r)i}}.$$
 (31)

# V. PERFORMANCE COMPUTATION OF WR-IM

The WR-IM performance is computed following the electromagnetic model and inductance calculations presented in the previous sections. First the estimated impedance matrix **Z** of (22) at an initial estimated current matrix **I** of (21) is determined. With **Z** known, a new (updated) **I** is determined from (23) at a set voltage **V** according to (20). For example, by only considering a positive q-axis voltage ( $\delta_s = 0$  of Fig. 2), the set **V** is given by

$$\begin{bmatrix} V_{ds} & V_{qs} & V_{dr} & V_{qr} \end{bmatrix}^T = \begin{bmatrix} 0 & 325 & 0 & 0 \end{bmatrix}^T.$$
 (32)

An iterative technique is then employed to update  $\mathbf{Z}$  at the updated  $\mathbf{I}$  until *rms* values of  $\mathbf{I}$  converge to a final value.

Fig. 3 shows the proposed calculation method (PCM) for the DQ-axis parameters and performance prediction with frozen



Fig. 4. IH-FE (left) and AM-FE (right) 2-D axial view models of the WR-IM.

permeabilities procedure. The variable  $\theta_m$  is the identified mechanical rotor position angle. In Fig. 3, k = 1, 2, ..., M is the convergence iterator where M is the number of iterations for the iterative technique termination. The current *rms* value of the  $k^{th}$  convergence iterator is calculated from the peak current value of (21) as

$$I_{rms_{(k)}} = \left\{ 0.5 \left[ I_{ds_{(k)}}^2 + I_{qs_{(k)}}^2 \right] \right\}^{\frac{1}{2}} \quad k = 1, 2, \dots, M.$$
(33)

The iterative technique is terminated when

$$I_{rms_{(k)}} - I_{rms_{(k-1)}} \leqslant \zeta \left(\frac{I_{rms_{(k)}} + I_{rms_{(k-1)}}}{2}\right) \tag{34}$$

where  $\zeta$  is the fraction tolerance in the current value and  $I_{rms_{(k)}}$  and  $I_{rms_{(k-1)}}$  are the start and final current values of one iteration. The DQ-axis current and torque values can also be used as function values for the iteration termination as in (34). Finally, note that the PCM of Fig. 3 for the performance computation of the WR-IM is done at a single rotor position, an aspect which is considered later in this paper.

#### VI. CASE STUDY COMPUTATION OF WR-IM

In this section, performance and parameter computed results are given of an unskewed WR-IM using the PCM implemented in an in-house static FEA (IH-FE) software package. The computed results are validated using a commercial transient FEA package, Ansoft/Maxwell (AM-FE). Fig. 4 shows 2-D IH-FE and AM-FE pole section simulation models of a four pole WR-IM with 36 stator and 24 rotor slots. The dimensions and parameters of the motor models of Fig. 4 are tabulated in Table I.

## A. Performance Computation Results

Fig. 5 shows (a) *rms* current values of (33)  $I_{(33)}^f$  and (b) torque values of (25)  $T_{(25)}^f$  versus the convergence iterator *k* of Fig. 3 for 15 iterations (i.e., M = 15). Different slip frequency values (i.e., f = 0, 1.5, and 50 Hz) shown on the superscript of the performance motor variables are used to demonstrate the PCM behavior at diverse conditions. At a set voltage matrix **V** of (32), the results here show a converging behavior which indicates an opportunity for a performance solution at **V** and *f*. It is shown that it takes only up to six iterations to obtain a solution to the motor performance independent of *f*.

#### **B.** Parameter Computation Results

This section presents parameter results versus rotor position using the PCM for the current solution in Fig. 5 at f =1.5 Hz (i.e.,  $n_r = 1$  455 r/min) and k = 6. Fig. 6 shows

TABLE I DIMENSION AND PARAMETER VALUES OF THE WR-IM

Dimension description	Stator	Rotor	Units
Outer diameter	210	144.4	(mm)
Inner diameter	145	40	(mm)
Slot opening height	1	1	(mm)
Slot wedge height	1.5	1.5	(mm)
Slot height	11.65	29.3	(mm)
Slot opening width	2	1	(mm)
Lower slot opening width	7.14	2.7	(mm)
Upper slot opening width	9.18	10.4	(mm)
Teeth width	5.95	8.61	(mm)
Stack length	120	120	(mm)
Air gap length	0.305		(mm)
Parameter description			
Phase resistance	1.26	0.73	(Ω)
End-winding inductance	0.01	0.01	(H)
Turns per slot	24	36	-
Rated frequency	50	48.5	(Hz)
Rated phase voltage Y connection	230	0	(V)
Rated phase current	6.8	5.9	(A)
Rated Power	3.4		(kW)



Fig. 5. (a) *rms* current and (b) torque values versus convergence iterator at slip frequencies of 0, 1.5 and 50 Hz.

the non-linear IH-FE solutions of total flux linkages of (17), also equal to the sum of the segregated flux linkages of (29) in (31). Fig. 7 shows the WR-IM self and mutual inductances values of (27) versus rotor position. These inductances are computed as of (30) from the segregated DQ-axis flux linkage components of (29). In Fig. 7(b), the results of only two mutual cross-coupling inductances are shown as an example. The results of  $M_{dsdr}$  and  $M_{qsqr}$  are not shown since they are equal to  $M_{drds}$  and  $M_{qrqs}$ , respectively. Fig. 8 shows the segregated torque components of (25)  $T_{(25)}$  (i.e.,  $T_t$ ) and the IH-FE torque  $T_{(IH)}$  versus rotor position.

The parameter and torque performance results of Figs. 7 and 8 with fixed DQ-axis currents show very little variation with rotor position for the unskewed motor. However,  $T_{(IH)}$ 



Fig. 6. Total flux linkages from the non-linear FE solution versus rotor position according to (9).



Fig. 7. DQ-axis (a) self and (b) mutual inductances versus rotor position.



Fig. 8. Torque versus rotor position.

has a high torque ripple as shown because it is the actual FE-calculated rotor position torque.

#### C. PCM Solution at Different Rotor Positions

As mentioned earlier, the results of Figs. 7 and 8 versus rotor position are with fixed DQ-currents. The question is how will the PCM current and torque solution of Fig. 3 vary with the designer's choice of rotor position,  $\theta_m$ . Fig. 9 shows these PCM results versus the designer's choice of  $\theta_m$ , with slip frequency a parameter. For the unskewed motor, the results show that the PCM predicts the current and torque at any choice of  $\theta_m$  without introducing a weighty error.



Fig. 9. rms current and torque versus rotor position with slip frequencies of (a) 0, (b) 1.5, and (c) 50 Hz.



Fig. 10. PCM and AM-FE rms current versus slip frequency.

# D. Validation of PCM Results

To validate the PCM of Fig. 3, the AM-FE model of the WR-IM of Fig. 4 is simulated. Fig. 10 shows the PCM and AM-FE stator and rotor *rms* currents (i.e.,  $I_{(s,r)(33)}$ ,  $I_{(s,r)(AM)}$ ) versus slip frequency, respectively. Fig. 11 also shows the IH-FE, and AM-FE average and PCM torques (i.e.,  $\overline{T}_{(IH)}$ ,  $\overline{T}_{(AM)}$ ,  $T_{(25)}$ ) versus slip frequency, respectively. The obtained results using the PCM are in excellent agreement with those obtained from the transient AM-FE and static IH-FE solutions.

# VII. DISCUSSION OF SIMULATION RESULTS

From the PCM of Fig. 3, a single iteration (M = 1) requires  $1 \times \text{non-linear plus } 4 \times \text{linear FE}$  solutions for parameter and performance prediction. Since the linear FE solution with



Fig. 11. IH-FE, AM-FE, and PCM torque versus slip frequency.

frozen permeabilities is relatively very fast, the time taken for the 4 × linear FE solutions can be neglected. Furthermore, for convergence from the results of Fig. 5, only M = 6 static non-linear FE solutions are needed for accurate parameter and performance prediction.

In order to have a quantitative comparison of the computational time, the steady state rms stator current solutions of the AM-FE and PCM at 50 Hz (stand still) and 1.5 Hz slip frequencies have been simulated. On a 2.93 GHz CPU with 12 GB RAM Intel(R) Core i7 computer, the execution times were found in this case to be 1833 and 166 s, respectively, for the AM-FE and PCM methods. Hence the PCM method is eleven times faster than the AM-FE method, which is extremely important when it comes to design optimization. Furthermore, the accessibility of the actual motor parameters (e.g., inductances) from the analysis using the total flux linkages gives the PCM a major advantage in the precise modeling of the electrical machine. Note has to be taken that any FEA package (commercial or non-commercial) can be utilized in the PCM as long as the package can execute static FE solutions and is capable of freezing the core-element permeabilities.

The PCM includes the effects of saturation saliency and cross-magnetization on the WR-IM circuit parameters [30]–[32]. However, it was found, among others from the results of Fig. 7, that the small WR-IM does not suffer significantly from the effect of cross-magnetization. Hence, the motor is easily analyzed on orthogonal axes assuming that they are decoupled. The inductance matrix of (18) in (17) can then classically be adapted as [2], [33]

$$\mathbf{L}_{(s,r)} \approx \begin{bmatrix} L_{ds} & 0 & M_d & 0\\ 0 & L_{qs} & 0 & M_q\\ M_d & 0 & L_{dr} & 0\\ 0 & M_q & 0 & L_{qr} \end{bmatrix}$$
(35)

where  $M_d = M_{drds} = M_{dsdr}$  and  $M_q = M_{qrqs} = M_{qsqr}$ . However, it is wrong to assume that the corresponding DQ-axis self and mutual inductances of the stator and rotor are equal as illustrated in (35). With the simplified inductance matrix of (35) the torque components of (26) in (25) simplifies to

$$T_r = \frac{3}{2} p(M_d I_{dr} I_{qs} - M_q I_{qr} I_{ds})$$
  

$$T_s = \frac{3}{2} p[(L_{ds} - L_{qs}) I_{ds} I_{qs}]$$
  

$$T_m = 0$$
(36)

as is also shown by the torque components in Fig. 8.

## VIII. CONCLUSION

An implementation method on modeling the WR-IM, essentially predicting motor parameters and performances at any slip speed, is described in this paper. The PCM is found to be successful in that it determines accurate and fast solutions for performance prediction. Hence, the method is extremely suited, amongst others, to the design optimization of the WR-IM.

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